

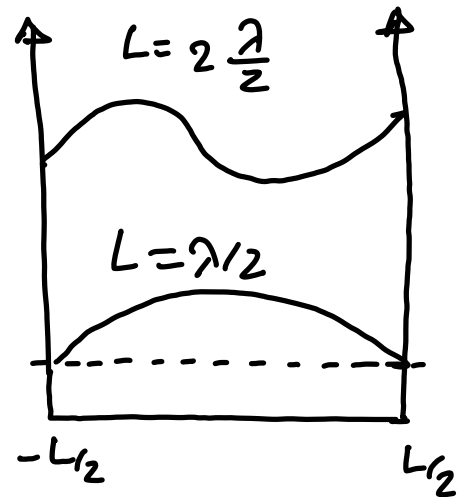
Infinite QW

$$\psi = A e^{ikx} + B e^{-ikx}$$

$$\boxed{L = n \frac{\lambda}{2}} \rightarrow \lambda = \frac{2L}{n}$$

$$k = \frac{2\pi}{\lambda} = \frac{2\pi}{\frac{2L}{n}} = \frac{\pi n}{L}$$

$$E = \frac{\hbar^2 k^2}{2m} \Rightarrow E_n = \frac{\hbar^2}{2m} \frac{n^2 \pi^2}{L^2}$$



Check B.C.:

$$A e^{-ikL/2} + B e^{ikL/2} = 0$$

$$\rightarrow A e^{-in\pi/2} + B e^{in\pi/2} = 0 \rightarrow$$

$$\begin{cases} n \text{ Even: } A(1) + B(1) = 0 \Rightarrow A = -B \\ n \text{ Odd: } A(-i) + B(i) = 0 \Rightarrow A = B \end{cases} \Rightarrow$$

$$\begin{cases} n \text{ Even: } \psi(x) = A (e^{-ikx} - e^{ikx}) = \underbrace{-2iA}_{C} \sin kx \\ n \text{ Odd: } \psi(x) = A' (e^{-ikx} + e^{ikx}) = \underbrace{2A'}_{C'} \cos kx \end{cases}$$

Eigenfunct. must be normalized:

$$c^2 \int_{-L/2}^{L/2} \sin^2 kx \, dx = c^2 \int_{-L/2}^{L/2} \sin^2 \frac{n\pi}{L} x \, dx = 1$$

$$\Rightarrow c^2 \left(\frac{L}{2} \right) = 1 \Rightarrow c = \sqrt{\frac{2}{L}}$$

Similarly for $c' = \sqrt{\frac{2}{L}} \Rightarrow$

$$\left\{ \begin{array}{l} \psi(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L} x \quad n \text{ even (odd parity)} \\ \psi(x) = \sqrt{\frac{2}{L}} \cos \frac{n\pi}{L} x \quad n \text{ odd (even parity)} \end{array} \right.$$

We can combine these in a single form:

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin \left(\frac{n\pi}{L} \left(x + \frac{L}{2} \right) \right) \quad \text{all } n$$

The complete wavefunction is:

$$\Psi_n(x, t) = \psi_n(x) e^{-i\omega_n t} \quad E_n = \hbar \omega_n$$

$$\Psi_n(x, t) = \sqrt{\frac{2}{L}} \sin \left(\frac{n\pi}{L} \left(x + \frac{L}{2} \right) \right) e^{-i\omega_n t}$$

Current density in QM

From Maxwell's equ. we know that:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \mathbf{J} \quad \rho(r, t), \mathbf{J}(r, t)$$

In QM: $n = |\psi|^2$ Particle density

$$\rho = en = e|\psi|^2$$

$$\frac{\partial \rho}{\partial t} = e \frac{\partial}{\partial t} (\psi^* \psi) = e \underbrace{\frac{\partial \psi^*}{\partial t} \psi}_{\text{is for c.c.}} + e \psi^* \underbrace{\frac{\partial \psi}{\partial t}}_{\text{is for c.c.}} \Rightarrow$$

- is for c.c.

$$\downarrow \frac{1}{-i\hbar} \left(-\frac{\hbar^2}{2m} \nabla^2 \psi^* + V\psi^* \right) \quad \frac{1}{i\hbar} \left(-\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi \right)$$

$$\frac{\partial \rho}{\partial t} = \frac{e}{i\hbar} \left[\frac{\hbar^2}{2m} (\psi \nabla^2 \psi^* - \psi^* \nabla^2 \psi) - \cancel{V\psi^* \psi} + \cancel{V\psi^* \psi} \right]$$

$$= \frac{e\hbar}{i2m} (\psi \nabla^2 \psi^* - \psi^* \nabla^2 \psi)$$

$$= \frac{-ie\hbar}{2m} \vec{\nabla} \cdot (\psi \vec{\nabla} \psi^* - \psi^* \vec{\nabla} \psi) = -\nabla \cdot \mathbf{J}$$

$$\Rightarrow \boxed{\vec{\mathbf{J}} = \frac{ie\hbar}{2m} (\psi \vec{\nabla} \psi^* - \psi^* \vec{\nabla} \psi)}$$